

# A completeness theorem for general relativity

Judit Madarász and Gergely Székely

Rényi Institute of Mathematics, Budapest

Logic, Relativity and Beyond, 2015

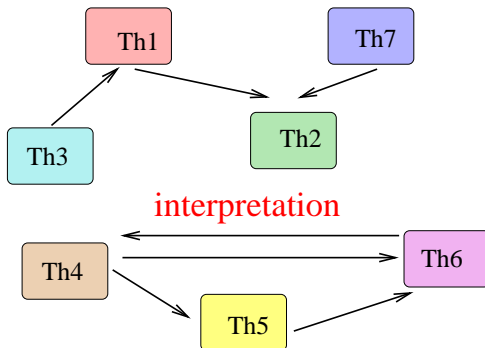
It was a dream of the Vienna Circle to build up scientific theories in mathematical logic



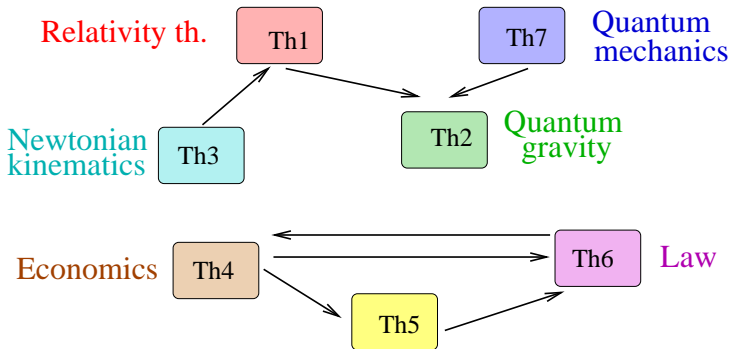
Benefits of building up a scientific theory (e.g. relativity theory) in mathematical logic:

- Uncover the tacit assumptions and make them explicit.
- Axiomatize theories (in the sense of math.logic).
- Analyze the relations between assumptions and consequences.
- Answer why type questions.
- Etc.

Breaking up a big theory into many smaller ones



## CONNECT DIFFERENT THEORIES, AREAS



# Relativity theory in first-order logic

## Andréka–Németi School/Team

Relativity theory in first-order logic

Andréka–Németi School/Team

Gergely Székely (after the coffee break) What structures can numbers have in relativity theory?

# Relativity theory in first-order logic

## Andréka–Németi School/Team

Gergely Székely (after the coffee break) What structures can numbers have in relativity theory?

Tommorow, from 4pm – 5:30pm

István Németi and Hajnal Andréka Relativity theory via a network of logic theories

Koen Lefever and Gergely Székely Interpretation of Special Relativity in the Language of Newtonian Kinematics

Attila Molnár Some Expressive Temporal Logic of Minkowski Spacetimes



# Relativity theory in first-order logic

## Andréka–Németi School/Team

Gergely Székely (after the coffee break) What structures can numbers have in relativity theory?

Tommorow, from 4pm – 5:30pm

István Németi and Hajnal Andréka Relativity theory via a network of logic theories

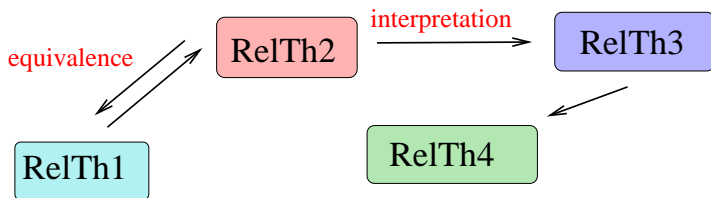
Koen Lefever and Gergely Székely Interpretation of Special Relativity in the Language of Newtonian Kinematics

Attila Molnár Some Expressive Temporal Logic of Minkowski Spacetimes

Wednesday 5pm

Mike Stannett Using an Automated Theorem Prover to Support First Order Relativity Theory

Tommorow, from 4pm – 5:30pm



## Axiomatization in general:

### Axioms:

Ax.1.

Ax.2.

Ax.3.

Etc.



### Theorems:

Thm.1.

Thm.2.

Thm.3.

Etc.

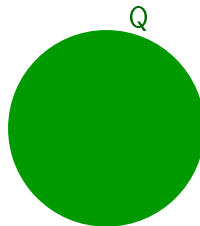
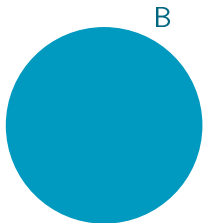


Economical  
Streamlined  
Transparent  
Observational



Rich  
Complex

Logic Language: { B, Q, }



B  $\leftrightarrow$  Bodies (things that move)

Q  $\leftrightarrow$  Quantities

Logic Language: { B, IOb, Ph, Q, }

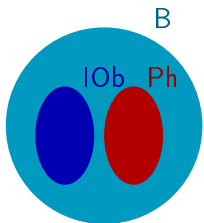


B  $\longleftrightarrow$  Bodies (things that move)

IOb  $\longleftrightarrow$  Inertial Observers      Ph  $\longleftrightarrow$  Photons (light signals)

Q  $\longleftrightarrow$  Quantities

Logic Language:  $\{ B, \text{IOb}, \text{Ph}, Q, +, \cdot, \leq, \}$



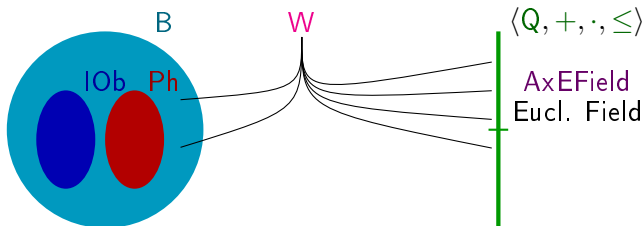
$\langle Q, +, \cdot, \leq \rangle$   
+ AxEField  
Eucl. Field

$B \iff$  Bodies (things that move)

$\text{IOb} \iff$  Inertial Observers      $\text{Ph} \iff$  Photons (light signals)

$Q \iff$  Quantities      $+, \cdot$  and  $\leq \iff$  field operations and ordering

Logic Language:  $\{ B, IOb, Ph, Q, +, \cdot, \leq, W \}$



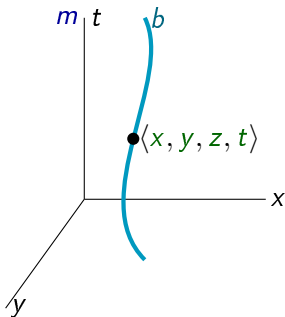
$B \leftrightarrow$  Bodies (things that move)

$IOb \leftrightarrow$  Inertial Observers     $Ph \leftrightarrow$  Photons (light signals)

$Q \leftrightarrow$  Quantities     $+, \cdot$  and  $\leq \leftrightarrow$  field operations and ordering

$W \leftrightarrow$  Worldview (a 6-ary relation of type  $BBQQQQ$ )

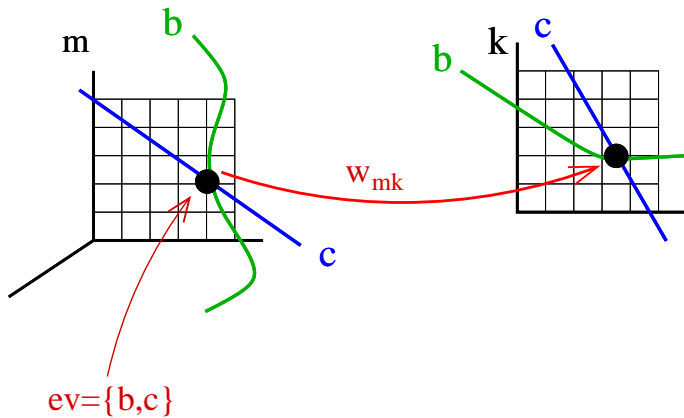
$W(m, b, x, y, z, t) \iff$  “observer  $m$  coordinatizes body  $b$  at spacetime location  $\langle x, y, z, t \rangle$ .”



Worldline of body  $b$  according to observer  $m$

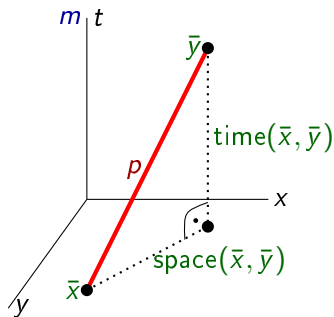
$$wline_m(b) = \{ \langle x, y, z, t \rangle \in Q^4 : W(m, b, x, y, z, t) \}$$





AxPh :

For any *inertial observer*, the *speed of light* is 1. Furthermore, it is possible to send out a *light signal* in any *direction*.



$$\begin{aligned} \forall m \big( \text{IOb}(m) \rightarrow \forall \bar{x} \bar{y} \big( \exists p [ \text{Ph}(p) \wedge \text{W}(m, p, \bar{x}) \wedge \text{W}(m, p, \bar{y}) ] \leftrightarrow \\ \leftrightarrow \text{space}(\bar{x}, \bar{y}) = \text{time}(\bar{x}, \bar{y}) \big) \end{aligned}$$

AxEv :

*Inertial observers coordinatize the same events.*

AxSelf :

*Every inertial observer is stationary according to himself.*

AxSym :

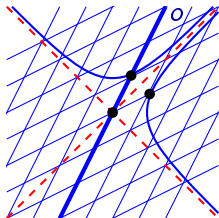
*Inertial observers use the same units of measurements.*

$$\text{SpecRel} = \text{AxPh} + \text{AxEField} + \text{AxEv} + \text{AxSelf} + \text{AxSym}$$

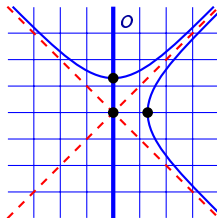
$$\text{SpecRel} = \text{AxPh} + \text{AxEField} + \text{AxEv} + \text{AxSelf} + \text{AxSym}$$

Theorem:

$\text{SpecRel} \models$  “The worldview transformations between inertial observers are Poincaré transformations.”



worldview of  $o'$

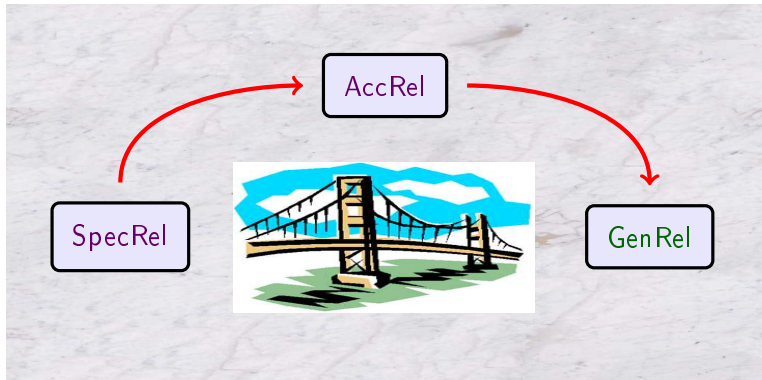


worldview of  $o$

Theorem:

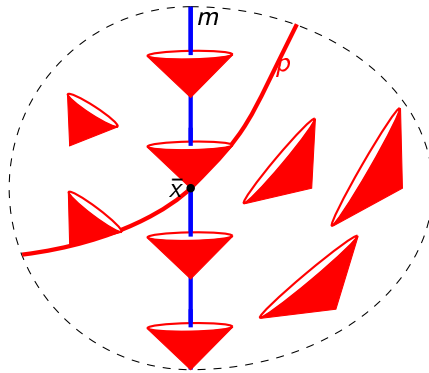
$\text{SpecRel} \Rightarrow \left\{ \begin{array}{l} \text{“Relatively moving clocks slow down.”} \\ \text{“Relatively moving spaceships shrink.”} \\ \text{etc.} \end{array} \right.$

# General Relativity



$AxPh^-$  :

If an *observer* sends out a *light signal*, then the speed of the *light-signal* is 1 according to the *observer*, and it is possible to send out a *light signal* in any *direction*.



# General Relativity

$$\text{GenRel} = \mathbf{AxPh}^- + \mathbf{AxSm}^- + \dots$$

Official Relativity Th

Not a FOL theory

?

GenRel

Our FOL General  
Relativity Th.

# General Relativity

$$\text{GenRel} = \mathbf{AxPh}^- + \mathbf{AxSm}^- + \dots$$

Official Relativity Th



LorMan



GenRel

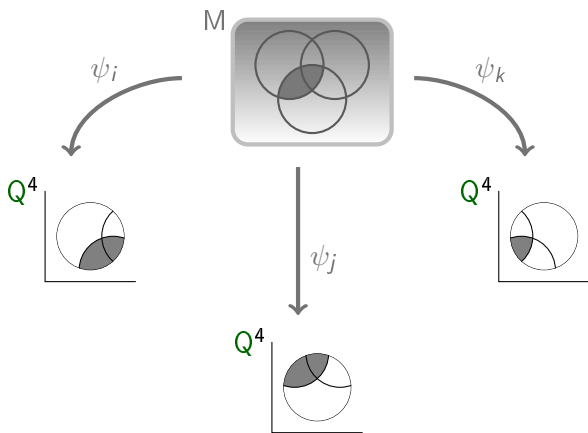
Our FOL General  
Relativity Th.

FOL theory of Loretzian manifolds



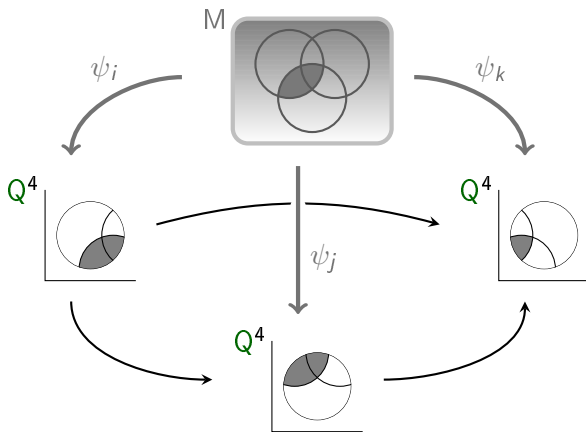
## Lorentzian Manifolds

$(M, g)$ ,  $M$  is a 4 dim. Man. and  $g_p : TM_p \times TM_p \rightarrow \mathbb{Q}$  Lorentzian metric



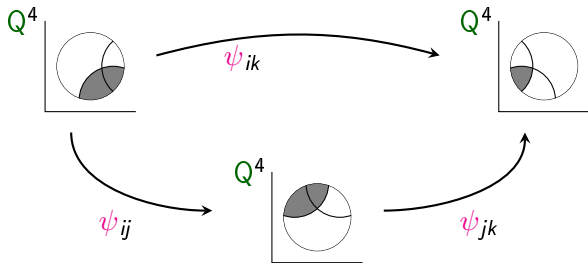
## Lorentzian Manifolds

$(M, g)$ ,  $M$  is a 4 dim. Man. and  $g_p : TM_p \times TM_p \rightarrow \mathbb{Q}$  Lorentzian metric



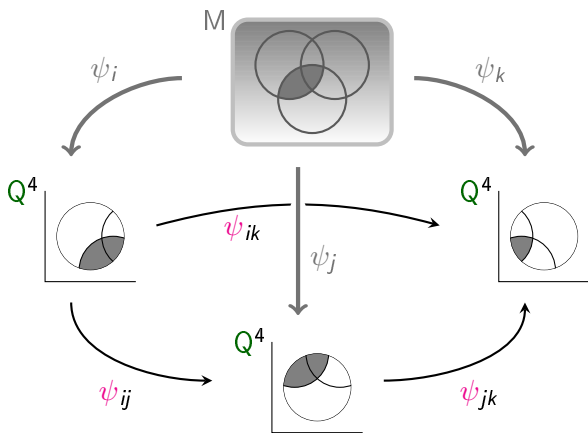
## Lorentzian Manifolds

$(M, g)$ ,  $M$  is a 4 dim. Man. and  $g_p : TM_p \times TM_p \rightarrow \mathbb{Q}$  Lorentzian metric



## Lorentzian Manifolds

$(M, g)$ ,  $M$  is a 4 dim. Man. and  $g_p : TM_p \times TM_p \rightarrow \mathbb{Q}$  Lorentzian metric



FOL language for Lorentzian manifolds

$\{ I, Q, \}$

$I \iff$  (Indexes) of charts     $Q \iff$  Quantities,

FOL language for Lorentzian manifolds

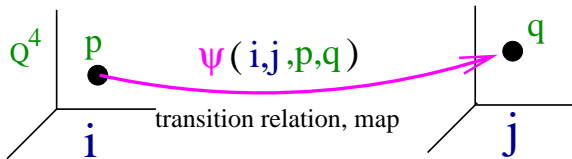
$\{I, Q, +, \cdot, \leq, \quad\}$

$I \longleftrightarrow$  (Indexes) of charts     $Q \longleftrightarrow$  Quantities,     $+, \cdot, \leq$

FOL language for Lorentzian manifolds

$$\{I, Q, +, \cdot, \leq, \psi, \}$$

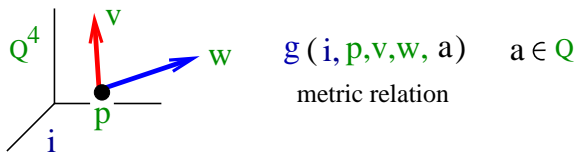
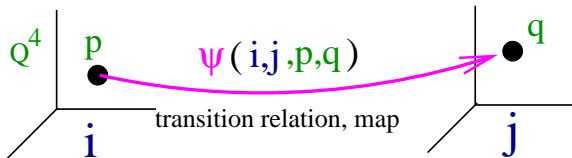
$I \leftrightarrow$  (Indexes) of charts     $Q \leftrightarrow$  Quantities,     $+, \cdot, \leq$



FOL language for Lorentzian manifolds

$$\{I, Q, +, \cdot, \leq, \psi, g\}$$

$I \leftrightarrow$  (Indexes) of charts     $Q \leftrightarrow$  Quantities,     $+, \cdot, \leq$





LorMan is a FOL theory of Lorentzian Manifolds.

**LorMan** is a FOL theory of Lorentzian Manifolds.

$$\mathcal{G} \models \text{GenRel}$$

$$\mathcal{G} \xrightarrow{TR} \mathcal{M} \models \text{LorMan}$$

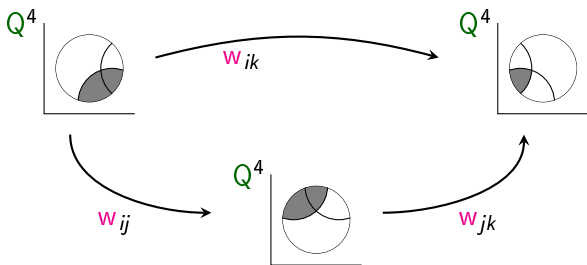
$$tr : Fm(\text{LorMan}) \longrightarrow Fm(\text{GenRel})$$

Theorem:

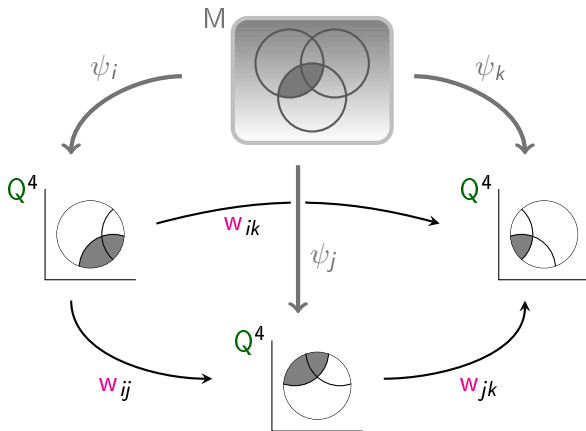
**GenRel** is complete with respect to the “standard models of GR”, i.e., Lorentzian manifolds.

$$\text{GenRel} \models tr(\text{LorMan})$$

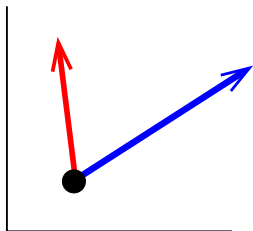
# GenRel Models

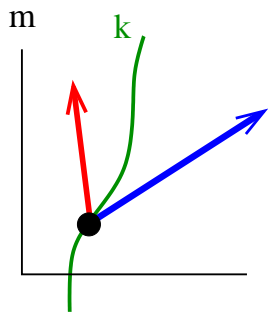


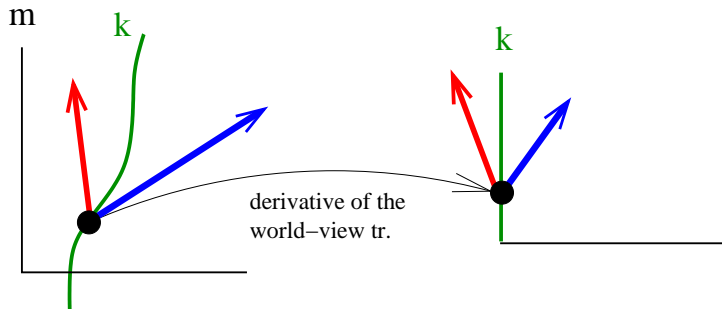
# GenRel Models

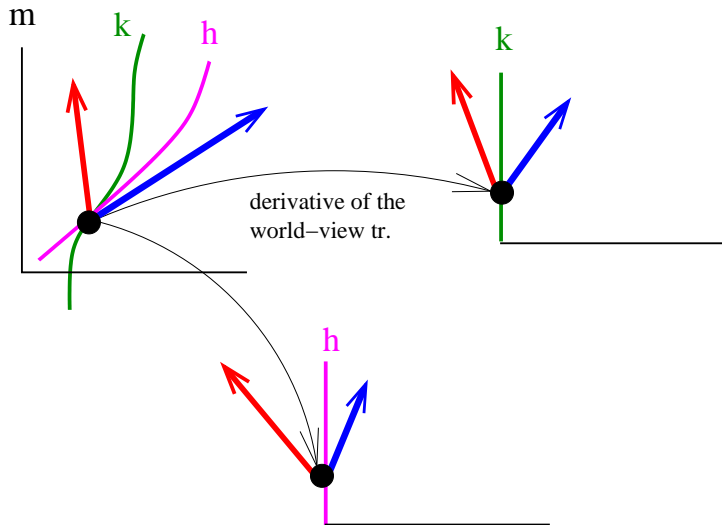


m

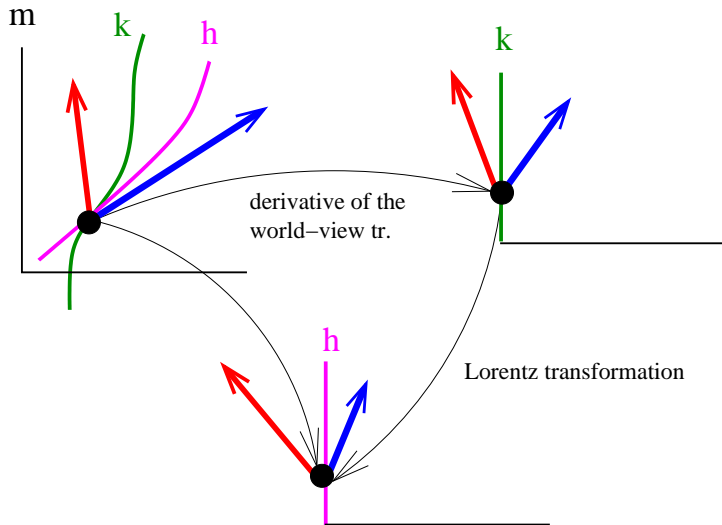




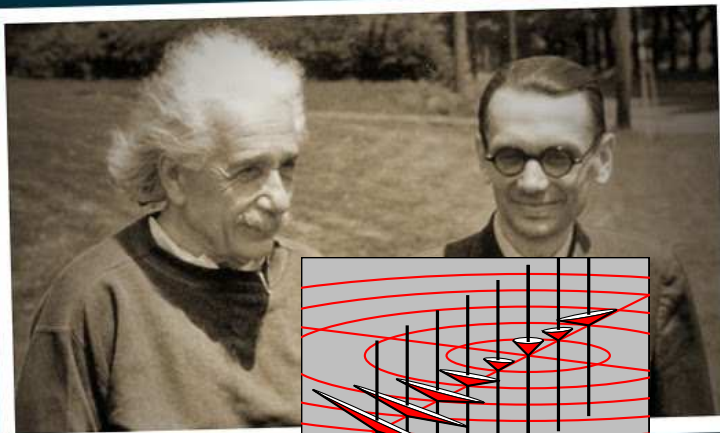








What about the other direction???



Einstein and Gödel (Princeton, 1948)

$\mathcal{M}$  is a Lorentzian Manifold over the Reals

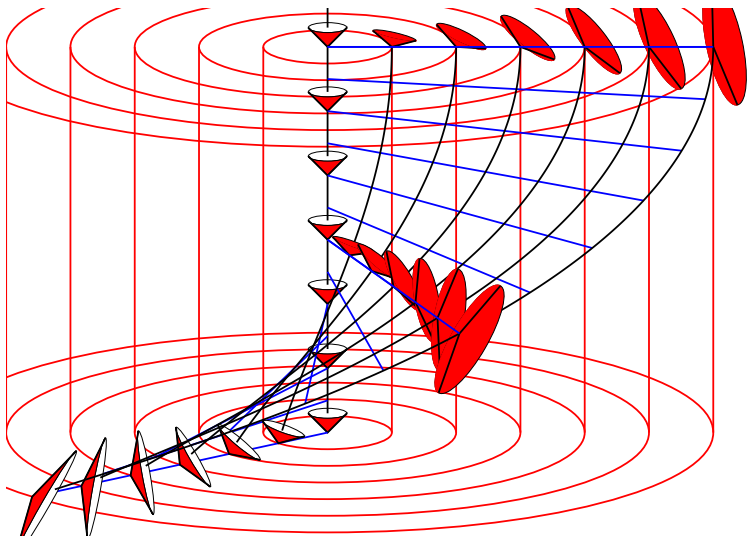
$$\mathcal{M} \xrightarrow{TR'} \mathcal{G} \models \text{GenRel}$$

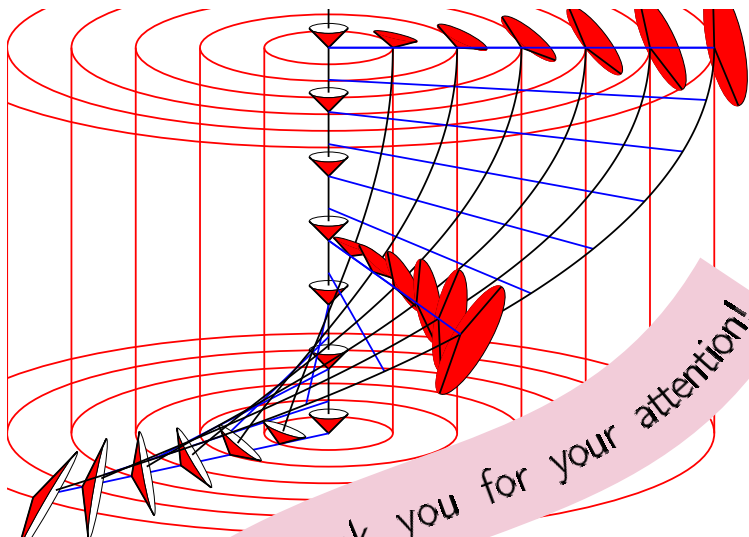
Theorem:

For any Lorentzian Manifold  $\mathcal{M}$  over the Reals, if

$$\mathcal{M} \xrightarrow{TR'} \mathcal{G} \xrightarrow{TR} \mathcal{N}$$

then  $\mathcal{M}$  and  $\mathcal{N}$  are isometric, thus they represent the same spacetime.





Thank you for your attention!