

# Beyond the spherical sup-norm problem

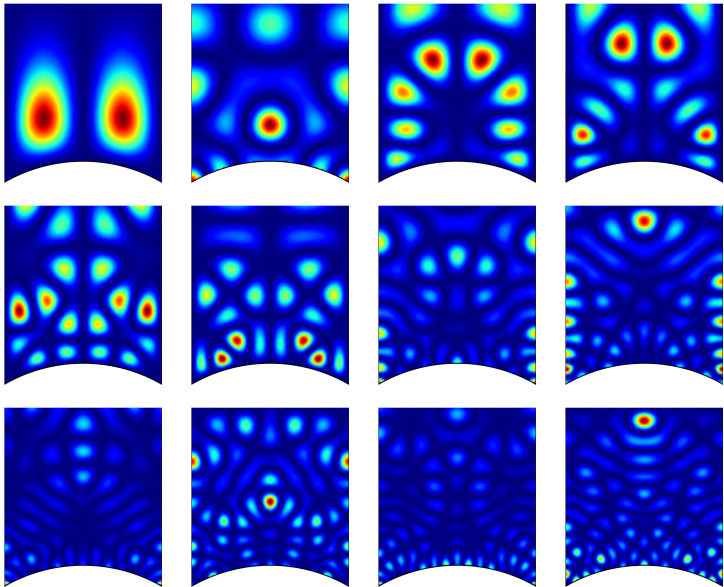
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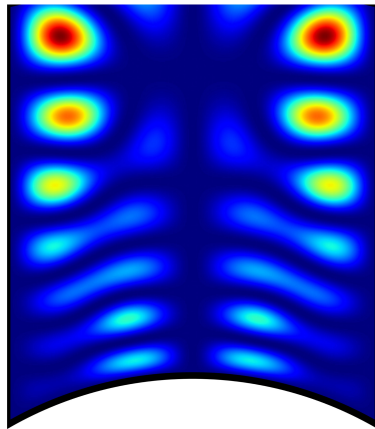
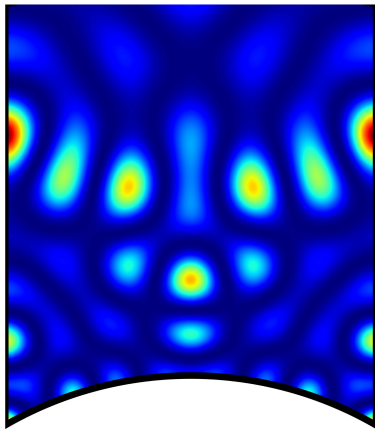
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ICTS, Bengaluru, India

# Maass forms on the modular surface

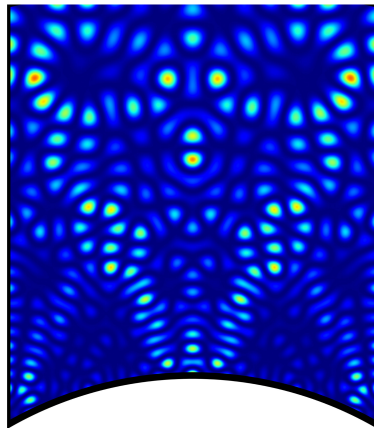
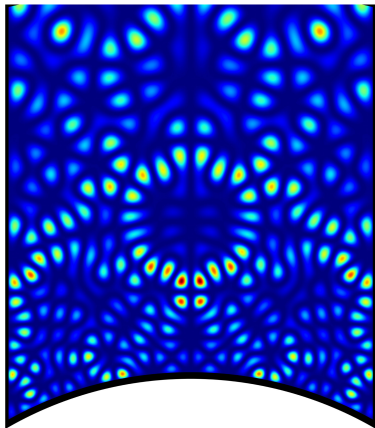


# Maass forms with $\lambda \approx 10^3$ on the modular surface



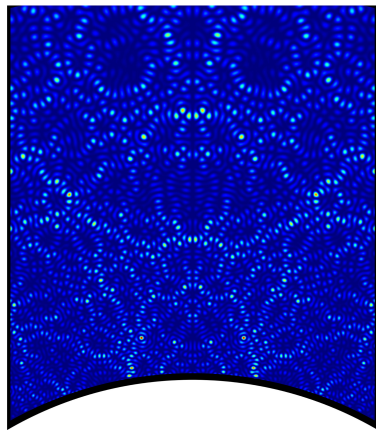
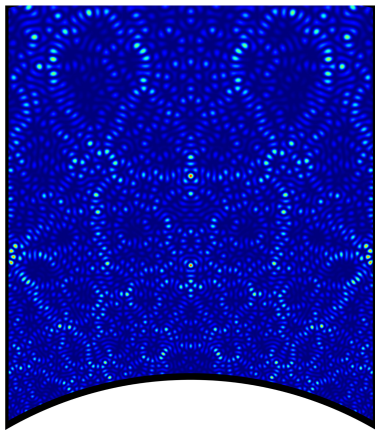
Fredrik Strömberg

# Maass forms with $\lambda \approx 10^4$ on the modular surface



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# Maass forms with $\lambda \approx 10^5$ on the modular surface



Fredrik Strömberg

# Classical mechanics vs. quantum mechanics

Let  $M$  be a compact orientable Riemannian manifold, and consider a particle moving freely on  $M$  with unit speed.

	classical mechanics	quantum mechanics
<b>phase space</b>	$SM$ sphere bundle	$L^2(M)$ Hilbert space
<b>moving particle</b>	$f : \mathbb{R} \rightarrow SM$ smooth	$\psi : \mathbb{R} \rightarrow L^2(M)$ $\ \psi\  = 1$
<b>bounded observable</b>	$a : SM \rightarrow \mathbb{R}$ smooth	$\text{Op}(a) : L^2(M) \rightarrow L^2(M)$ self-adjoint & bounded
<b>time evolution</b>	$G^t : SM \rightarrow SM$ geodesic flow	$U_t : L^2(M) \rightarrow L^2(M)$ $U_t = e^{-it\sqrt{\Delta}}$

## Solutions of the Schrödinger equation

$$\psi(t) = U_t(\psi(0)) = \sum_{j=0}^{\infty} c_j e^{-it\sqrt{\lambda_j}} \phi_j, \quad (c_j)_{j=0}^{\infty} \in \ell^2(\mathbb{N})$$

# Quantum ergodicity on the modular surface (1 of 2)

Theorem (Shnirelman 1974, Colin de Verdière 1985, Zelditch 1987)

Assume that the geodesic flow on  $SM$  is ergodic, and let  $\{\phi_j\}$  be an orthonormal basis of  $L^2(M)$  satisfying  $\Delta\phi_j = \lambda_j\phi_j$ . Consider  $d\omega_j$  defined via  $\langle \text{Op}(a)\phi_j, \phi_j \rangle = \int_{SM} a d\omega_j$  for  $a \in C^\infty(SM)$ . Then  $d\omega_j \xrightarrow{*} d\omega$  along a subsequence of  $\lambda_j$ 's of density 1.

Proof (sketch).

Assume that  $a \in C^\infty(SM)$  has space average  $\int_{SM} a d\omega = 0$ .

Consider also a fixed time average  $a^T := \frac{1}{T} \int_0^T a \circ G^t dt$ .

By Egorov, Cauchy–Schwarz, Weyl, and Birkhoff, we have

$$\begin{aligned} \frac{1}{N(\lambda, 1)} \sum_{\lambda_j \leq \lambda} \left| \int_{SM} a d\omega_j \right|^2 &= \frac{1}{N(\lambda, 1)} \sum_{\lambda_j \leq \lambda} \left| \int_{SM} a^T d\omega_j \right|^2 + o(1) \\ &\leq \frac{1}{N(\lambda, 1)} \sum_{\lambda_j \leq \lambda} \int_{SM} |a^T|^2 d\omega_j + o(1) = \int_{SM} |a^T|^2 d\omega + o(1) < \varepsilon, \end{aligned}$$

for  $T = T_0(\varepsilon)$  and  $\lambda > \lambda_0(\varepsilon)$ . Hence the left hand side is  $o(1)$ .  $\square$

# Quantum ergodicity on the modular surface (2 of 2)

## Theorem (Hopf 1936)

Let  $M := \mathrm{SL}_2(\mathbb{Z}) \backslash \mathcal{H}^2$  be the modular surface. The geodesic flow on the sphere bundle  $SM \cong \mathrm{SL}_2(\mathbb{Z}) \backslash \mathrm{SL}_2(\mathbb{R})$  is ergodic.

## Proof (sketch).

Assume that  $f \in L^2(\mathrm{SL}_2(\mathbb{Z}) \backslash \mathrm{SL}_2(\mathbb{R}))$  is fixed by the right action of positive diagonal matrices  $\begin{pmatrix} a & \\ & a^{-1} \end{pmatrix}$ .

Then, for any fixed  $b \in \mathbb{R}$  and for  $a > 0$  tending to infinity,

$$\begin{aligned} \left\| \begin{pmatrix} 1 & b \\ & 1 \end{pmatrix} f - f \right\| &= \left\| \begin{pmatrix} 1 & b \\ & 1 \end{pmatrix} \begin{pmatrix} a & \\ & a^{-1} \end{pmatrix} f - \begin{pmatrix} a & \\ & a^{-1} \end{pmatrix} f \right\| \\ &= \left\| \begin{pmatrix} a^{-1} & \\ & a \end{pmatrix} \begin{pmatrix} 1 & b \\ & 1 \end{pmatrix} \begin{pmatrix} a & \\ & a^{-1} \end{pmatrix} f - f \right\| \\ &= \left\| \begin{pmatrix} 1 & a^{-1}b \\ & 1 \end{pmatrix} f - f \right\| \rightarrow \|f - f\| = 0. \end{aligned}$$

Hence any upper triangular matrix in  $\mathrm{SL}_2(\mathbb{R})$  fixes  $f$ . Similarly, any lower triangular matrix in  $\mathrm{SL}_2(\mathbb{R})$  fixes  $f$ . In the end, the entire group  $\mathrm{SL}_2(\mathbb{R})$  fixes  $f$ , and so  $f$  is constant almost everywhere.  $\square$



# The spherical sup-norm problem (1 of 2)

Theorem (Iwaniec–Sarnak 1995)

Let  $\phi$  be an  $L^2$ -normalized Hecke–Maass form on  $SL_2(\mathbb{Z}) \backslash \mathcal{H}^2$  of Laplacian eigenvalue  $\lambda$ . Then for any  $\varepsilon > 0$  we have

$$\|\phi\|_\infty \ll_\varepsilon \lambda^{5/24+\varepsilon}.$$

Using that  $\mathcal{H}^2 \cong SL_2(\mathbb{R})/SO_2(\mathbb{R})$ , the functions  $\phi$  above can be thought of as functions on  $SL_2(\mathbb{Z}) \backslash SL_2(\mathbb{R})/SO_2(\mathbb{R})$ .

They span the subspace of  $L^2_{\text{cusp}}(SL_2(\mathbb{Z}) \backslash SL_2(\mathbb{R}))$  fixed by the right action of  $SO_2(\mathbb{R})$ . That is, they span the **spherical** subspace of the cuspidal subspace.

In adelic language, the functions  $\phi$  fixed by  $T_{-1}$  live on

$$PGL_2(\mathbb{Q}) \backslash PGL_2(\mathbb{A}_{\mathbb{Q}}) / PO_2(\mathbb{R}) \prod_p PGL_2(\mathbb{Z}_p).$$

So these “even” functions  $\phi$  are **spherical** at every place of  $\mathbb{Q}$ .

# The spherical sup-norm problem (2 of 2)

Theorem (Blomer–Harcos–Milićević 2016)

Let  $\phi$  be an  $L^2$ -normalized Hecke–Maass form on  $SL_2(\mathbb{Z}[i]) \backslash \mathcal{H}^3$  of Laplacian eigenvalue  $\lambda$ . Then for any  $\varepsilon > 0$  we have

$$\|\phi\|_\infty \ll_\varepsilon \lambda^{5/12+\varepsilon}.$$

Using that  $\mathcal{H}^3 \cong SL_2(\mathbb{C})/SU_2(\mathbb{C})$ , the functions  $\phi$  above can be thought of as functions on  $SL_2(\mathbb{Z}[i]) \backslash SL_2(\mathbb{C})/SU_2(\mathbb{C})$ .

They span the subspace of  $L^2_{\text{cusp}}(SL_2(\mathbb{Z}[i]) \backslash SL_2(\mathbb{C}))$  fixed by the right action of  $SU_2(\mathbb{C})$ . That is, they span the **spherical** subspace of the cuspidal subspace.

In adelic language, the functions  $\phi$  fixed by  $T_i$  live on

$$PGL_2(\mathbb{Q}(i)) \backslash PGL_2(\mathbb{A}_{\mathbb{Q}(i)}) / PU_2(\mathbb{C}) \prod_{\mathfrak{p}} PGL_2(\mathbb{Z}[i]_{\mathfrak{p}}).$$

So these “even” functions  $\phi$  are **spherical** at every place of  $\mathbb{Q}(i)$ .

# Spectral decomposition of the automorphic $L^2$ space

$\Gamma$	$G$	$K$
$SL_2(\mathbb{Z})$	$SL_2(\mathbb{R})$	$SO_2(\mathbb{R})$
$SL_2(\mathbb{Z}[i])$	$SL_2(\mathbb{C})$	$SU_2(\mathbb{C})$
$PGL_2(\mathbb{Q})$	$PGL_2(\mathbb{A}_{\mathbb{Q}})$	$PO_2(\mathbb{R}) \prod_p PGL_2(\mathbb{Z}_p)$
$PGL_2(\mathbb{Q}(i))$	$PGL_2(\mathbb{A}_{\mathbb{Q}(i)})$	$PU_2(\mathbb{C}) \prod_p PGL_2(\mathbb{Z}[i]_p)$

We have the following Hilbert space decompositions into **irreducible  $G$ -spaces** and their **spherical (i.e.  $K$ -invariant)** subspaces:

$$L^2_{\text{cusp}}(\Gamma \backslash G) = \bigoplus_{\pi} V_{\pi},$$

$$L^2_{\text{cusp}}(\Gamma \backslash G / K) = L^2_{\text{cusp}}(\Gamma \backslash G)^K = \bigoplus_{\pi} V_{\pi}^K.$$

For the adelic groups, the above decompositions are unique (multiplicity one), and the spaces  $V_{\pi}$  consist of Hecke eigenforms. The nonzero subspaces  $V_{\pi}^K$  are the one-dimensional spaces  $\mathbb{C}\phi$ , where  $\phi$  is a **spherical** Hecke–Maass form as before.

# The unitary dual of $G = \mathrm{SL}_2(\mathbb{R})$

The unitary dual  $\widehat{G}$  was determined by [Bargman \(1947\)](#), perhaps inspired by the work of [Wigner \(1939\)](#). The nontrivial irreducible unitary representations of  $G$  are infinite dimensional, and the ones relevant here (the tempered ones) come in 4 families:

- **spherical/non-spherical principal series**  $\pi_{it}^\pm$  for  $t \in \mathbb{R}_{\geq 0}$ ;
- **holomorphic/antiholomorphic discrete series**  $\pi_k^\pm$  for  $k \in \mathbb{Z}_{\geq 1}$ .

These representations can be defined explicitly, e.g. by letting  $G$  act on  $L^2(\mathbb{R})$  in natural but different ways. We can clearly distinguish between the above 4 types by looking at how they decompose into **irreducible  $K$ -spaces** (we parametrize  $\widehat{K}$  by  $\mathbb{Z}$ ):

$$\begin{aligned} V_{it}^+ &= \bigoplus_{\substack{\ell \in \mathbb{Z} \\ \ell \equiv 0 \pmod{2}}} V_{it}^{+, \ell} & V_{it}^- &= \bigoplus_{\substack{\ell \in \mathbb{Z} \\ \ell \equiv 1 \pmod{2}}} V_{it}^{-, \ell} \\ V_k^+ &= \bigoplus_{\substack{\ell \geq k \\ \ell \equiv k \pmod{2}}} V_k^{+, \ell} & V_k^- &= \bigoplus_{\substack{\ell \leq -k \\ \ell \equiv k \pmod{2}}} V_k^{-, \ell} \end{aligned}$$

The summands here are one-dimensional (i.e. isomorphic to  $\mathbb{C}$ ).

# The non-spherical sup-norm problem for $SL_2(\mathbb{Z}) \backslash SL_2(\mathbb{R})$

Inspired by the theory of newforms in the level aspect, the natural **non-spherical sup-norm problem** would concern a minimal weight vector in  $\pi_k^+$  (or equivalently a maximal weight vector in  $\pi_k^-$ ), which is again unique up to scaling. However, such a vector has weight  $k$  Laplacian eigenvalue  $\frac{k}{2} (1 - \frac{k}{2})$ , which grows with  $k$ . So the weight aspect is not separated from the eigenvalue aspect in this variant of the problem.

To go genuinely beyond the spherical sup-norm problem, we are led to work with  $SL_2(\mathbb{C})$  rather than  $SL_2(\mathbb{R})$ .

If  $\pi_k^+$  occurs in  $L_{\text{cusp}}^2(\Gamma \backslash G)$ , then *the absolute value* of its minimal weight vector is invariant under  $K$ , and as a function on  $\mathcal{H}^2$  it agrees with  $F(x + iy) := y^{k/2} |f(x + iy)|$ , where  $f$  is a holomorphic cusp form of weight  $k$  and level 1 on  $\mathcal{H}^2$ . If  $f$  is a Hecke eigenform and  $\|F\|_2 = 1$ , then  $\|F\|_\infty \ll_\varepsilon k^{1/4+\varepsilon}$  by a result of [Xia \(2007\)](#). For co-compact  $\Gamma \leq SL_2(\mathbb{R})$ , a similar result was proved by [Khayutin–Steiner \(2020\)](#), improving on [Das–Sengupta \(2015\)](#).

# The unitary dual of $G = \mathrm{SL}_2(\mathbb{C})$

The unitary dual  $\widehat{G}$  was determined by **Gelfand–Naimark (1947)**. The nontrivial irreducible unitary representations of  $G$  are infinite dimensional, and the ones relevant for the moment (the tempered ones) come in a single family:

- **principal series**  $\pi_{it,p}$  for  $t \in \mathbb{R}_{\geq 0}$  and  $p \in \frac{1}{2}\mathbb{Z}$ .

These representations can be defined explicitly, e.g. by letting  $G$  act on  $L^2(\mathbb{C})$  in a natural way.

It is instructive (and crucial for us) to look at how these representations decompose into **irreducible  $K$ -spaces** (we parametrize  $\widehat{K}$  by  $\frac{1}{2}\mathbb{Z}_{\geq 0}$ ), and further into one-dimensional subspaces under the action of the diagonal subgroup of  $K$ :

$$V_{it,p} = \bigoplus_{\substack{\ell \geq |p| \\ \ell \equiv p \pmod{1}}} \bigoplus_{\substack{\ell \geq |p| \\ \ell \equiv p \pmod{1}}} \bigoplus_{\substack{|q| \leq \ell \\ q \equiv \ell \pmod{1}}} V_{it,p}^{\ell,q}.$$

Here  $\dim V_{it,p}^{\ell} = 2\ell + 1$  and  $\dim V_{it,p}^{\ell,q} = 1$ .

## Notation

$$\Gamma := \mathrm{SL}_2(\mathbb{Z}[i]), \quad G := \mathrm{SL}_2(\mathbb{C}), \quad K := \mathrm{SU}_2(\mathbb{C}).$$

## Theorem (Blomer–Harcos–Maga–Milićević 2021)

Let  $\ell \geq 1$  be an integer,  $I \subset \mathbb{R}$  and  $\Omega \subset G$  be compact sets. Let  $V_\pi \subset L^2_{\mathrm{cusp}}(\Gamma \backslash G)$  be a cuspidal representation such that  $\pi \simeq \pi_{it,p}$ , where  $t \in I$  and  $|p| = \ell$ . As usual, we assume that  $V_\pi$  consists of Hecke eigenfunctions. Let us choose an orthonormal basis  $\{\phi_q : |q| \leq \ell\}$  of  $V_\pi^\ell$ , with  $\phi_q \in V_\pi^{\ell,q}$ . Then for any  $\varepsilon > 0$  we have

$$\sum_{|q| \leq \ell} |\phi_q(g)|^2 \ll_{\varepsilon, I, \Omega} \ell^{8/3 + \varepsilon}, \quad g \in \Omega.$$

For the individual summands we have

$$\phi_q(g) \ll_{\varepsilon, I, \Omega} \ell^{26/27 + \varepsilon}, \quad g \in \Omega.$$

Finally, for  $q = 0$  (resp. for  $q = \pm \ell$  under a technical assumption), we can improve the exponent to  $7/8 + \varepsilon$  (resp. to  $1/2 + \varepsilon$ ).

Following [Selberg \(1956\)](#), consider a rapidly decaying continuous function  $f \in L^1(G)$ , and its action on  $L^2(\Gamma \backslash G)$  given by

$$(R(f)\psi)(g) := \int_G f(h)\psi(gh) \, dh = \int_{\Gamma \backslash G} \left( \sum_{\gamma \in \Gamma} f(g^{-1}\gamma h) \right) \psi(h) \, dh.$$

Assume that  $R(f)$  is a [positive operator](#), and  $\pi(f)$  acts by a scalar  $c(\pi, \ell)$  on  $V_\pi^\ell$ . Then  $R(f)$  preserves the orthogonal decomposition  $V_\pi^\ell \oplus V_\pi^{\ell, \perp}$ . Moreover,  $R(f)$  composed with the projection to  $V_\pi^\ell$  has a simple kernel just like  $R(f)$ :

$$(R(f)_\pi^\ell \psi)(g) = \int_{\Gamma \backslash G} \left( c(\pi, \ell) \sum_{|q| \leq \ell} \phi_q(g) \overline{\phi_q(h)} \right) \psi(h) \, dh.$$



## Bounding the sup-norm via an automorphic kernel (2 of 2)

By a simple approximation argument, on the diagonal  $g = h$ , the kernel of  $R(f)_\pi^\ell$  is upper bounded by the kernel of  $R(f)$ :

$$c(\pi, \ell) \sum_{|q| \leq \ell} |\phi_q(g)|^2 \leq \sum_{\gamma \in \Gamma} f(g^{-1}\gamma g), \quad g \in G.$$

To make this work in practice, we consider all functions  $f \in L^2(G)$  for which every  $\pi_{it,p}(f)$  acts by a scalar on the component  $V_{it,p}^\ell$ , and by zero on the other components  $V_{it,p}^m$ . These functions form a Hilbert subspace  $\mathcal{H}(\tau_\ell) \subset L^2(G)$  defined by the conditions

- $f(g) = f(kgk^{-1})$  for almost every  $g \in G$  and  $k \in K$ ;
- $f = \bar{\chi}_\ell \star f \star \bar{\chi}_\ell$ , where  $\chi_\ell$  is  $2\ell + 1$  times the character of  $\tau_\ell$ .

For  $f \in L^1(G) \cap \mathcal{H}(\tau_\ell)$  the scalar  $c(\pi, \ell)$  exists, and for  $\pi \cong \pi_{it,p}$  it equals  $\hat{f}(it, p)/(2\ell + 1)$ , where

$$\hat{f}(it, p) := \text{tr}(\pi_{it,p}(f)) = \int_G f(g) \varphi_{it,p}^\ell(g) dg,$$
$$\varphi_{it,p}^\ell(g) := \text{tr}(\pi_{it,p}(\bar{\chi}_\ell) \pi_{it,p}(g) \pi_{it,p}(\bar{\chi}_\ell)).$$

# Generalized spherical transform

The theory of **Gelfand–Naimark (1947 & 1950)** yields the Hilbert space isomorphism  $\mathcal{H}(\tau_\ell) \cong L^2(\widehat{G}(\tau_\ell))$  with the **Plancherel identity**

$$\int_G |f(g)|^2 dg = \frac{1}{2\ell + 1} \sum_{\substack{|\rho| \leq \ell \\ \rho \equiv \ell \pmod{1}}} \int_0^\infty |\widehat{f}(it, \rho)|^2 (t^2 + \rho^2) dt.$$

In practice we define  $f(g)$  in terms of its **generalized spherical transform**  $\widehat{f}(it, \rho)$  using the **inversion formula**

$$f(g) = \frac{1}{2\ell + 1} \sum_{\substack{|\rho| \leq \ell \\ \rho \equiv \ell \pmod{1}}} \int_0^\infty \widehat{f}(it, \rho) \varphi_{it, \rho}^\ell(g^{-1}) (t^2 + \rho^2) dt.$$

We need to ensure that  $f(g)$  is continuous, rapidly decaying, and of reasonable size. For this we need to understand the **spherical trace function**  $\varphi_{it, \rho}^\ell(g)$  in some detail.

# Spherical trace function

The spherical trace function has an integral representation over  $K$  involving the diagonal matrix coefficients of  $\tau_\ell$ . As a result, it extends holomorphically to  $(\nu, p, g) \in \mathbb{C} \times \frac{1}{2}\mathbb{Z} \times G$ , and it satisfies a soft general bound that we skip for simplicity. In particular,

$$\varphi_{\nu,\ell}^\ell(g) = (2\ell + 1) \int_K \kappa_\ell(k^{-1}gk) dk,$$

where

$$\kappa_\ell \left( \begin{pmatrix} a & b \\ c & d \end{pmatrix} \right) := \bar{a}^{2\ell} (|a|^2 + |c|^2)^{\nu-\ell-1}.$$

The spherical trace function also has remarkable symmetries:

$$\varphi_{\nu,p}^\ell(g) = \overline{\varphi_{-\bar{\nu},p}^\ell(g)} = \varphi_{\nu,p}^\ell(g^{-1}),$$

$$\varphi_{\nu,p}^\ell(g) = \varphi_{p,\nu}^\ell(g), \quad \nu \equiv p \pmod{1}, \quad |\nu|, |p| \leq \ell.$$

These properties become transparent by analytically continuing the representations  $\pi_{it,p}$  to (non-unitary Fréchet) representation  $\pi_{\nu,p}$ .

# Paley–Wiener space and Schwartz space

Now we see that if  $f \in L^1(G) \cap \mathcal{H}(\tau_\ell)$  decays rapidly, then its transform  $\widehat{f}$  extends holomorphically to  $\mathbb{C} \times \frac{1}{2}\mathbb{Z}$  such that

$$\widehat{f}(\nu, p) = \widehat{f}(p, \nu), \quad \nu \equiv p \pmod{1}, \quad |\nu|, |p| \leq \ell. \quad (*)$$

This means that  $\nu$  and  $p$  are not independent as we thought!

## Theorem (Wang 1974)

For  $f \in \mathcal{H}(\tau_\ell)$  and  $R > 0$ , the following conditions are equivalent.

- 1  $f$  is smooth, and  $f(k_1 a_h k_2) = 0$  for  $|h| > R$  and  $k_1, k_2 \in K$ .
- 2  $\widehat{f}$  extends holomorphically to  $\mathbb{C} \times \frac{1}{2}\mathbb{Z}$  such that  $(*)$  holds true, and we also have  $\widehat{f}(\nu, p) \ll_C (1 + |\nu|)^{-C} e^{R|\Re \nu|}$ .

## Theorem (Blomer–Harcos–Maga–Milićević 2021)

For  $f \in \mathcal{H}(\tau_\ell)$ , the following conditions are equivalent.

- 1  $f$  is smooth, and  $\frac{\partial^m}{\partial h^m} f(k_1 a_h k_2) \ll_{m,A} e^{-A|h|}$  for  $k_1, k_2 \in K$ .
- 2  $\widehat{f}$  extends holomorphically to  $\mathbb{C} \times \frac{1}{2}\mathbb{Z}$  such that  $(*)$  holds true, and we also have  $\widehat{f}(\nu, p) \ll_{B,C} (1 + |\nu|)^{-C}$  for  $|\Re \nu| \leq B$ .

# Choice of test function

We ended up using the function  $f \in \mathcal{H}(\tau_\ell)$  whose transform equals

$$\widehat{f}(\nu, \rho) = \begin{cases} e^{(p^2 - \ell^2 + \nu^2)/2}, & \nu \in \mathbb{C}, \quad p \in \frac{1}{2}\mathbb{Z}, \quad |p| \leq \ell; \\ 0, & \nu \in \mathbb{C}, \quad p \in \frac{1}{2}\mathbb{Z}, \quad |p| > \ell. \end{cases}$$

This provides a positive operator  $R(f)$  on  $L^2(\Gamma \backslash G)$  such that

$$c(\pi, \ell) \gg 1/\ell \quad \text{and} \quad f(g) \ll \ell^2 e^{-\log^2 \|g\|}.$$

However, this only yields the **baseline bound**

$$\sum_{|q| \leq \ell} |\phi_q(g)|^2 \ll \ell^3.$$

In order to improve on this, we need to **amplify**  $\pi$  by Hecke operators. This idea was introduced by **Iwaniec–Sarnak (1995)**.

# Amplification (1 of 2)

Using a standard amplifier, we can bound the sum of  $|\phi_q(g)|^2$  by

$$L^{-2+\varepsilon} \ell^2 \sum_{\substack{\gamma \in M_2(\mathbb{Z}[i]) \\ n = \det \gamma \neq 0 \\ \|g^{-1} \tilde{\gamma} g\| \leq \ell^\varepsilon}} \frac{|x_n|}{|n|} \sup_{\nu \in i\mathbb{R}} |\varphi_{\nu, \ell}^\ell(g^{-1} \tilde{\gamma} g)| + L^{2+\varepsilon} \ell^{-48},$$

where  $\tilde{\gamma}$  abbreviates  $\gamma / \sqrt{\det \gamma}$  for any choice of  $\sqrt{\det \gamma}$ , and

- $x_1 \ll L$ ;
- $x_n \ll 1$  when  $n$  equals  $l_1 l_2$  or  $l_1^2 l_2^2$  for two split Gaussian primes  $l_1, l_2$  of length about  $\sqrt{L}$  and angle in  $(0, \pi/4)$ ;
- $x_n = 0$  otherwise.

## Theorem (Blomer–Harcos–Maga–Milićević 2021)

Let  $\ell \geq 1$  be an integer, and let  $h = \begin{pmatrix} z & u \\ & z^{-1} \end{pmatrix} \in G$  be upper triangular. Then for any  $\nu \in i\mathbb{R}$ ,  $k \in K$ ,  $\varepsilon > 0$ , we have

$$\varphi_{\nu, \ell}^\ell(k^{-1} h k) \ll_\varepsilon \min \left( \ell, \frac{\ell^\varepsilon \|h\|^6}{|z^2 - 1|^2}, \frac{\ell^{1/2+\varepsilon} \|h\|^3}{|u|} \right).$$

## Amplification (2 of 2)

We defined  $f$  in terms of the spherical trace function  $\varphi_{\nu,p}^{\ell}$  using the inversion formula. Replacing  $\varphi_{\nu,p}^{\ell}$  by

$$\varphi_{\nu,p}^{\ell,q}(h) := \frac{1}{2\pi} \int_0^{2\pi} \varphi_{\nu,p}^{\ell}(h \operatorname{diag}(e^{i\varrho}, e^{-i\varrho})) e^{-2qi\varrho} d\varrho$$

in this definition has the effect of picking a single  $\phi_q$ :

$$|\phi_q(g)|^2 \leq L^{-2+\varepsilon} \ell^2 \sum_{\substack{\gamma \in M_2(\mathbb{Z}[i]) \\ n = \det \gamma \neq 0 \\ \|g^{-1}\tilde{\gamma}g\| \leq \ell^\varepsilon}} \frac{|x_n|}{|n|} \sup_{\nu \in i\mathbb{R}} |\varphi_{\nu,\ell}^{\ell,q}(g^{-1}\tilde{\gamma}g)| + L^{2+\varepsilon} \ell^{-48}.$$

**Theorem (Blomer–Harcos–Maga–Milićević 2021)**

Let  $\ell, q \in \mathbb{Z}$  be such that  $\ell \geq \max(1, |q|)$ . Let  $\nu \in i\mathbb{R}$  and  $h \in G$ . Then for any  $\varepsilon > 0$  and  $\Lambda > 0$ , we have

$$\varphi_{\nu,\ell}^{\ell,q}(h) \ll_{\varepsilon,\Lambda} \ell^\varepsilon \min \left( 1, \frac{\|h\|}{\sqrt{\ell} \operatorname{dist}(h, K)^2 \operatorname{dist}(h, \mathcal{D})} \right) + \ell^{-\Lambda}.$$