

Constructibility and Space-Time



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Principal Results

- ❑ Space-time is developed from a constructible set theoretical foundation
- ❑ The real line is countable and functions of a real variable are locally homeomorphic with the real line
- ❑ Space-time is relational and its differential properties fulfill the requirements of Einstein-Weyl causality
- ❑ The Schrödinger equation is developed in this theory
- ❑ This result also infers that quantum mechanics in this relational space-time framework can be considered conceptually cumulative with prior physics.

Agenda

- (1) Axioms of a constructible set theory are presented.
- (2) Continuous and differentiable functions of a real variable are developed in this theory.
- (3) A nonlinear sigma model is postulated and the Schrödinger equation is shown to be a special case.
- (4) Implications of the theory to space-time are then presented.
- (5) Other examples are given of questions in physics that now may be answered by this approach.

ZF – Axiom Schema of Subsets – Power Set + Constructibility

Extensionality	Two sets with just the same members are equal.
Pairs	For every two sets, there is a set that contains just them.
Union	For every set of sets, there is a set with just all their members.
Infinity	There is at least one set ω^* which contains the null set and for every member there is a next member that contains just all its predecessors
Arithmetic	The Peano axioms without the induction axiom
Regularity	Every non-empty set has a minimal member (i.e. “weak” regularity).
Replacement	Replacing members of a set one-for-one creates a set. (i.e. bijective replacement)
Constructibility	Subsets of ω^* are countably constructible.

The power set axiom and the axiom schema of subsets are deleted and an axiom of constructibility adjoined.

Then:

- The theory is uniformly dependent on ω^*
- Both finite and infinite natural numbers exist in ω^*
- The set of constructible subsets of ω^* is countable
- All sets of finite natural numbers are finite
- Convergence of “infinite series” is undecidable.

The set of all ratios of any two elements of ω^* is called Q^* .

Q^* is an enlargement of the usual set of rational numbers Q .

Two members of Q^* are called “identical” if their ratio is 1.

We can now employ the symbol “ $\equiv$ ” for “is identical to”.

An “infinitesimal” is a member of Q^* “equal” to 0, i.e., letting y signify the member and employing the symbol “=” to signify equality, $y = 0 \leftrightarrow \forall n[y < 1/n]$ where n is an integer.

The reciprocal of an infinitesimal is “infinite”. A member of Q^* that is not an infinitesimal and not infinite is “finite”.

Note that $x \equiv y \rightarrow x = y$ but not the converse.

The constructibility axiom well-orders the set of constructible subsets of ω^* , creating a metric space. These subsets are binimals making up a real line $\mathbf{R}^*$.

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An *equality-preserving* bijective map $\Phi(x,u)$ between intervals X and U of $\mathbf{R}^*$ in which $x \in X$ and $u \in U$

$$\forall x_1, x_2, u_1, u_2 \left[\Phi(x_1, u_1) \wedge \Phi(x_2, u_2) \rightarrow [x_1 - x_2 = 0 \leftrightarrow u_1 - u_2 = 0] \right]$$

creates pieces biunique and homeomorphic with $\mathbf{R}^*$.

Note that U can be infinitesimal if and only if X is, i.e., the range vanishes if and only if the domain vanishes.

$u(x)$ is a function of a real variable in this theory if and only if it is a constant or a sequence in x of continuously connected biunique pieces such that the derivative with respect to x is also a function of a real variable. These functions are of bounded variation, locally homeomorphic with the real line $\mathbf{R}^*$ and, if not constant, the range of $u(x)$ is finite ($\text{range } u(x) \neq 0$) if and only if its domain is finite.

If some derivative is a constant, these are polynomials. If no derivative is a constant, these functions do not exist in this theory. They can, however, be approached as closely as necessary by a sum of polynomials of sufficiently high degree obtained by many iterations of the following algorithm, thereby creating effectively a finite dimensional Hilbert space:

$$\int_a^b \left[p \left(\frac{du}{dx} \right)^2 - qu^2 \right] dx \equiv \lambda; \quad \int_a^b ru^2 dx \equiv \text{Const.}$$

where $a \neq b$; $u \frac{du}{dx} \equiv 0$ at a and b ; p , q and r are functions of x .

Polynomials of sufficiently high degree are obtained by iteratively minimizing λ for a given p , q and r . They are effectively Sturm-Liouville eigenfunctions and are of bounded variation and locally homeomorphic with $\mathbb{R}^*$.

This has both physical and philosophical implications.

Consider a fixed string of finite length:

$$\int \left[\left(\frac{\partial u_x u_t}{\partial x} \right)^2 - \left(\frac{\partial u_x u_t}{\partial t} \right)^2 \right] dx dt \equiv 0$$

The u_x and u_t can be effectively obtained by iteration subject to the constraint $\lambda_x - \lambda_t \equiv 0$

This can be generalized to more complex fields and to finitely many dimensions of a compactified space-time.

Let $\mathbf{u}_{\ell mi}(\mathbf{x}_i)$ and $\mathbf{u}_{\ell mj}(\mathbf{x}_j)$ be eigenfunctions with positive eigenvalues $\lambda_{\ell mi}$ and $\lambda_{\ell mj}$ respectively.

We now define a “field” as a sum of eigenstates

$$\underline{\Psi}_m = \sum_{\ell} \Psi_{\ell m} \underline{i}_{\ell}, \Psi_{\ell m} = \prod_i \mathbf{u}_{\ell mi} \prod_j \mathbf{u}_{\ell mj}$$

subject to the novel postulate that for every eigenstate m the value of the integral of the Lagrange density over $d\mathbf{s}d\tau$, where

$d\mathbf{s}d\tau = \prod_i r_i d\mathbf{x}_i \prod_j r_j d\mathbf{x}_j$, is *identically null*: For all m

$$\sum_{\ell} \int \left\{ \sum_i \frac{1}{r_i} \left[P_{\ell mi} \left(\frac{\partial \Psi_{\ell m}}{\partial \mathbf{x}_i} \right)^2 - Q_{\ell mi} \Psi_{\ell m}^2 \right] - \sum_j \frac{1}{r_j} \left[P_{\ell mj} \left(\frac{\partial \Psi_{\ell m}}{\partial \mathbf{x}_j} \right)^2 - Q_{\ell mj} \Psi_{\ell m}^2 \right] \right\} d\mathbf{s}d\tau = 0$$

over a compactified space-time; this is a **nonlinear sigma model**

The postulate asserts that the two symmetric components of the nonlinear sigma model have identical magnitudes.

In this nonlinear sigma model the P and Q can be functions of any of the χ_i and χ_j , thus of any $\Psi_{\ell m}$ as well, subject to the condition that all the eigenvalues are finite and positive.

The Ψ_m are effectively given by iteration,
subject to the constraint

$$\sum_{\ell} \left(\sum_i \lambda_{\ell mi} - \sum_j \lambda_{\ell mj} \right) \equiv 0$$

Let both $\sum_m \sum_\ell \int \left\{ \sum_i \frac{1}{r_i} \left[P_{\ell mi} \left(\frac{\partial \Psi_{\ell m}}{\partial x_i} \right)^2 - Q_{\ell mi} \Psi_{\ell m}^2 \right] \right\} ds d\tau$ and

$$\sum_m \sum_\ell \int \left\{ \sum_j \frac{1}{r_j} \left[P_{\ell mj} \left(\frac{\partial \Psi_{\ell m}}{\partial x_j} \right)^2 - Q_{\ell mj} \Psi_{\ell m}^2 \right] \right\} ds d\tau \quad \text{be represented by } \alpha(\Psi)$$

- I. $\alpha(\Psi)$ is positive and closed to addition and to the absolute value of subtraction, so it is either continuous or discrete.
- II. $\alpha(\Psi) \equiv 0 \leftrightarrow \Psi \equiv 0$; otherwise, as Ψ is a function of real variables, the range $\Psi \neq 0$ and $\forall \Psi \alpha(\Psi) \neq 0$, thus $\alpha(\Psi)$ is not continuous.
- III. $\therefore \alpha(\Psi)$ is discrete; $\alpha(\Psi) \equiv n\kappa$, where n is any integer and κ is some finite unit which must be obtained empirically.

Let $\ell = 1, 2$, $r_\ell = P_{1\ell t} = P_{2\ell t} = 1$, $Q_{1\ell t} = Q_{2\ell t} = 0$, $\tau = \omega_m t$ and we normalize Ψ as follows:

$$\Psi_m = \sqrt{(C/2\pi)} \prod_i u_{im}(x_i) [u_{1m}(\tau) + i \cdot u_{2m}(\tau)] \quad (9)$$

where $i = \sqrt{-1}$ with

$$\int \sum_m \prod_i u_{im}^2 ds (u_{1m}^2 + u_{2m}^2) \equiv 1 \quad (10)$$

then:

$$\frac{du_{1m}}{d\tau} = -u_{2m} \quad \text{and} \quad \frac{du_{2m}}{d\tau} = u_{1m} \quad (11)$$

or

$$\frac{du_{1m}}{d\tau} = u_{2m} \quad \text{and} \quad \frac{du_{2m}}{d\tau} = -u_{1m} \quad (12)$$

For the minimal non-vanishing field, α is just the least finite value κ . Thus,

$$\begin{aligned} (C/2\pi) \sum_m \oint \int \left[\left(\frac{du_{1m}}{d\tau} \right)^2 + \left(\frac{du_{2m}}{d\tau} \right)^2 \right] \\ \prod_i u_{im}^2(x_i) ds d\tau \equiv C \equiv \kappa \end{aligned} \quad (13)$$

Substituting the Planck constant h for κ , this can now be put into the conventional Lagrangian form for the time term in the Schrödinger equation,

$$\frac{h}{2i} \sum_m \oint \int \left[\Psi_m^* \left(\frac{\partial \Psi_m}{\partial t} \right) - \left(\frac{\partial \Psi_m^*}{\partial t} \right) \Psi_m \right] ds dt$$

Going back to the statement $\forall \Psi \alpha(\Psi) \neq 0$, we must recognize we have assumed that space-time exists, i.e., that the upper and lower limits of at least one of the multiple integrals over all space-time are not equal. Otherwise, space-time does not exist $\rightarrow \forall \Psi \alpha(\Psi) = 0$ thus we obtain a necessary and sufficient condition:

$$\text{space-time exists} \leftrightarrow \exists \Psi \alpha(\Psi) \neq 0$$

This is a description of a relational space-time.

Furthermore, it has been shown by others that Q^2 is an ordered space that fulfills the strict Hausdorff topology requirements for Einstein-Weyl causality. Also, if Q^2 can be embedded in R^2 so that fields are locally homeomorphic with R , this will even apply to R^n (n is any positive integer).

In this theory, Q^{*2} is an ordered space one-to-one with R^{*2} , fields are by definition locally homeomorphic with R^* and so R^{*n} fulfills the requirements for Einstein-Weyl causality.

For Further Discussion

- ❑ This theory proposes that the physical universe is composed entirely of constructible sets; accordingly, the universe will not have non-constructible sets which can lead to physical antinomies. This has to be intuitively satisfying since, were there any physical antinomies, the universe would tear itself apart.
- ❑ Dyson's well-known problem, that the QED perturbation series cannot converge and thus diverges, does not apply in this theory. Instead, convergence of the perturbation series is undecidable. This provides a solution and opportunity for new physical ideas.
- ❑ Wigner's philosophical question concerning the unreasonable effectiveness of mathematics in physics may be answered directly; this effectiveness comes from their foundations being linked.

References

This is a partial list:

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ZF – Axiom Schema of Subsets – Power Set + Constructibility = T

Extensionality - Two sets with just the same members are equal

$$\forall x \forall y [\forall z [z \in x \leftrightarrow z \in y] \rightarrow x = y]$$

Pairs - For every two sets, there is a set that contains just them.

$$\forall x \forall y \exists z [\forall w [w \in z \leftrightarrow w = x \vee w = y]]$$

Union - For every set of sets, there is a set with just all their members.

$$\forall x \exists y \forall z [z \in y \leftrightarrow \exists u [z \in u \wedge u \in x]]$$

Infinity - A set ω^* contains 0 and every member has a successor containing just all its predecessors.

$$\exists \omega^* [0 \in \omega^* \wedge \forall x [x \in \omega^* \rightarrow x \cup \{x\} \in \omega^*]]$$

Replacement - Replacing members of a set one-for-one creates a set (i.e., “bijective” replacement).

Let $\phi(x,y)$ a formula in which x and y are free,

$$\forall z \forall x \in z \exists y [\phi(x,y) \wedge \forall u \in z \forall v [\phi(u,v) \rightarrow u = x \leftrightarrow y = v]] \rightarrow \exists r \forall t [t \in r \leftrightarrow \exists s \in z \phi(s,t)]$$

Regularity - Every non-empty set has a minimal member (i.e. “weak” regularity).

$$\forall x [\exists y y \in x \rightarrow \exists y [y \in x \wedge \forall z \neg [z \in x \wedge z \in y]]]$$

Constructibility - The subsets of ω^* are countably constructible.

$$\forall \omega^* \exists S [(\omega^*, 0) \in S \wedge \forall y \forall z [(y, z) \in S \rightarrow \\ [[\exists m_y [m_y \in y \wedge \forall v \neg [v \in y \wedge v \in m_y]] \leftrightarrow [\exists t_y \forall u [u \in t_y \leftrightarrow u \in y \wedge u \neq m_y]] \wedge (t_y \cup m_y, z \cup \{z\}) \in S]]].$$

Arithmetic:

The four formulae (a) to (d) below (Peano arithmetic axioms) can be consistently added to T.

We use x' for $x \cup \{x\}$.

$$\text{a) } \forall a \ a' \neq 0$$

$$\text{b) } \forall x \forall y \ (x' = y' \rightarrow x = y)$$

Let x, y, a, b be members of ω^* and $[x, y]$ and $[[x, y], z]$ represent ordered pairs.

$$\text{c) } \exists A \forall x \forall y \ \exists! z \ [[x, 0], x] \in A \wedge [[x, y], z] \in A \rightarrow [[x, y'], z'] \in A; \text{ this is addition: } x + y = z$$

$$\text{d) } \exists M \forall x \forall y \ \exists! z \ [[x, 0], 0] \in M \wedge [[x, y], z] \in M \rightarrow [[x, y'], z + x] \in M; \text{ this is multiplication: } x \cdot y = z$$

Define $[a, b]_r$ such that $[a_1, b]_r + [a_2, b]_r \equiv [a_1 + a_2, b]_r$ and $[a_1, b_1]_r \equiv [a_2, b_2]_r \leftrightarrow a_1 \cdot b_2 \equiv a_2 \cdot b_1$.

The extended set of rationals Q^* is the set of such pairs for all a and b in ω^* .

Theorem

$\forall \omega^* \forall x (\forall y_1 \forall y_2 \forall z_1 \forall z_2 \ ([x, y_1], z_1] \in B \wedge [[x, y_2], z_2] \in B \rightarrow ((z_1 = z_2) \leftrightarrow (y_1 = y_2))) \rightarrow$
 $\exists t(x) \ \forall z (z \in t(x) \leftrightarrow \exists y (B(x, y, z)))$, where B represents A or M .