# Interpretation of Special Relativity in the Language of Newtonian Kinematics

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Logic, Relativity & Beyond - Budapest, 2015.08.11

A translation is a function between formulas of languages preserving the logical connectives, i.e.  $Tr(\phi \wedge \psi) = Tr(\phi) \wedge Tr(\psi)$ , etc.

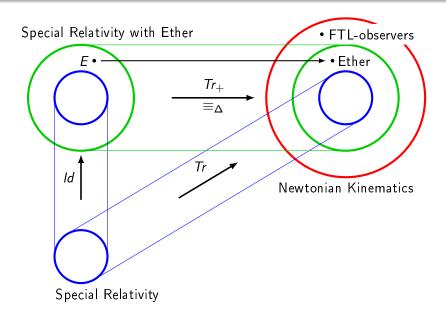
An interpretation of theory  $Th_1$  in theory  $Th_2$  is a translation Tr which translates all axioms of  $Th_1$  into theorems of  $Th_2$ .

An interpretation is a definitional equivalence if a translated formula can be translated back such that it becomes a formula which is equivalent to the original formula.

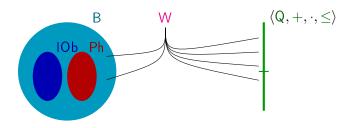
We use these concepts to show the *differences* between theories which are *not* equivalent.

There are translations Tr,  $Tr_+$ , and  $Tr_+^{-1}$  between the languages of NK and SR such that:

- $NK \vdash Tr(SR)$
- $NKnoFTL \vdash Tr_{+}(SRwithEther)$
- $SRwithEther \vdash Tr_{+}^{-1}(NKnoFTL)$
- Definitional equivalence:  $SRwithEther \equiv_{\Delta} NKnoFTL$ , i.e.,  $Tr_{+}$  and  $Tr_{+}^{-1}$  are inverses of each other up to logical equivalence in NKnoFTL and SRwithEther.



Language:  $\{B, IOb, Ph, Q, +, \cdot, \leq, W\}$ 



B ←→ Bodies (things that move)

IOb ←→ Inertial Observers Ph ←→ Photons (light signals)

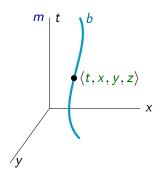
Q ↔ Quantities

+,  $\cdot$  and  $\leq \leftrightarrow \Rightarrow$  field operations and ordering

W ← Worldview (a 6-ary relation of type BBQQQQ)



 $W(m, b, t, x, y, z) \Leftrightarrow$  "observer m coordinatizes body b at spacetime location  $\langle t, x, y, z \rangle$ ."



Worldline of body b according to observer m

$$Wl_m(b):=\{\langle t,x,y,z\rangle\in\mathbb{Q}^4:\mathbb{W}(m,b,t,x,y,z)\}$$

 $\textbf{NewtonianKin}_{\textit{Full}} := \textbf{Kin} \cup \{ \texttt{AxEther}, \texttt{AbsTime}, \texttt{AxThExp}^{\uparrow}_{+} \}$ 

 $\textbf{SpecRel}_{\textit{Full}} := \textbf{Kin} \cup \{ \texttt{AxPh}_c, \texttt{AxThExp}^{\uparrow} \}$ 

## AxEField:

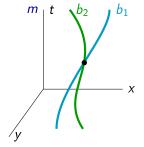
The structure of quantities  $(Q, +, \cdot, \leq)$  is an Euclidean field,

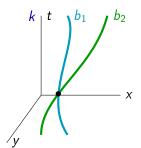
- Real numbers:  $\mathbb{R}$ ,
- Real algebraic numbers:  $\overline{\mathbb{Q}} \cap \mathbb{R}$ ,
- ullet Hyperreal numbers:  $\mathbb{R}^*$ ,
- Real constructable numbers,
- Etc...

# AxEv:

Inertial observers coordinatize the same events (meetings of bodies).

$$ev_m(\bar{x}) := \{b : W(m, b, \bar{x})\}$$



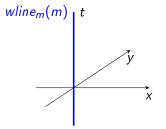


$$\forall m \, k\bar{x} \, [\mathsf{IOb}(m) \land \mathsf{IOb}(k) \rightarrow \exists \bar{y} \, ev_m(\bar{x}) = ev_k(\bar{y})].$$



## AxSelf:

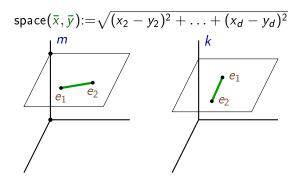
Every Inertial observer is stationary according to himself.



$$\forall mtxyz \ \Big( \mathsf{IOb}(m) \to \big[ \mathbf{W}(m,m,t,x,y,z) \leftrightarrow x = y = z = 0 \big] \Big).$$

# AxSymD :

Inertial observers agree as to the spatial distance between two events if these two events are simultaneous for both of them.



$$\forall mk\bar{x}\bar{y}\bar{x}'\bar{y}' \ [IOb(m) \land IOb(k) \land x_1 = y_1 \land x_1' = y_1' \land ev_m(\bar{x}) = ev_k(\bar{x}') \\ \land ev_m(\bar{y}) = ev_k(\bar{y}') \rightarrow \operatorname{space}(\bar{x}, \bar{y}) = \operatorname{space}(\bar{x}', \bar{y}')]$$

### AxLine:

The worldlines of inertial observers are straight lines according to inertial observers.

$$\forall mk\bar{x}\bar{y}\bar{z} \ (IOb(m) \land IOb(k) \land W(m,k,\bar{x}) \land W(m,k,\bar{y}) \land W(m,k,\bar{z})$$
$$\rightarrow \exists a \ [\bar{z} - \bar{x} = a(\bar{y} - \bar{x}) \lor \bar{y} - \bar{z} = a(\bar{z} - \bar{x})]).$$

(In SpecRel, AxLine is a theorem, so including it as an axiom is redundant.)

### AxTriv:

Any trivial transformation of an inertial frame is also an inertial frame.

 $\forall T \in Triv[\forall m \exists k(w_{mk} = T)]$ , where Triv is the set of Trivial transformations, i.e. transformations that are isometries on space and translations on time.

## AxAbsTime:

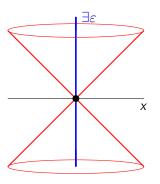
The time difference between two events is the same for all inertial observers.

$$\mathsf{time}(\bar{x},\bar{y}) := |x_1 - y_1|$$

$$\forall mk\bar{x}\bar{y}\bar{x}'\bar{y}' \ [IOb(m) \land IOb(k) \land ev_m(\bar{x}) = ev_k(\bar{x}') \land ev_m(\bar{y}) = ev_k(\bar{y}') \\ \rightarrow \mathsf{time}(\bar{x},\bar{y}) = \mathsf{time}(\bar{x}',\bar{y}')].$$

# AxEther(Einstein's AxLight):

There exists an inertial observer in which the light cones are right.

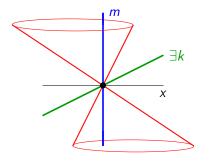


$$\exists \varepsilon c \Big[ \mathsf{IOb}(\varepsilon) \land c > 0 \land \forall \bar{x} \bar{y} \Big( \exists p \Big[ \mathsf{Ph}(p) \land \mathsf{W}(\varepsilon, p, \bar{x}) \Big] \\ \land \mathsf{W}(\varepsilon, p, \bar{y}) \Big] \leftrightarrow \mathsf{space}(\bar{x}, \bar{y}) = c \cdot \mathsf{time}(\bar{x}, \bar{y}) \Big) \Big]$$

# $AxThExp^{\uparrow}_{+}$ :

Inertial observers can move along any non-horizontal straight line.

$$k \uparrow m \stackrel{def}{\iff} ev_k(\bar{x}) = ev_m(1,0,0,0) \land ev_k(\bar{y}) = ev_m(0,0,0,0) \rightarrow x_1 > y_1$$

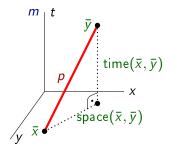


$$\exists h \ (IOb(h)) \land \forall m\bar{x}\bar{y} \ (IOb(m) \land x_1 \neq y_1 \\ \rightarrow \exists k \ IOb(k) \land W(m, k, \bar{x}) \land W(m, k, \bar{y}) \land m \uparrow k).$$

## Three pretactions Language Axit

# AxPhc:

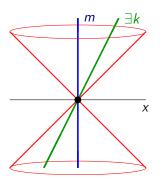
For any inertial observer, the speed of light is the same in every direction everywhere, and it is finite. Furthermore, it is possible to send out a light signal in any direction.



$$\exists c \Big[ c > 0 \land \forall m \bar{x} \bar{y} \Big( \mathsf{IOb}(m) \to \exists p \Big[ \mathsf{Ph}(p) \land \mathsf{W}(m, p, \bar{x}) \Big) \\ \land \mathsf{W}(m, p, \bar{y}) \Big] \leftrightarrow \mathsf{space}(\bar{x}, \bar{y}) = c \cdot \mathsf{time}(\bar{x}, \bar{y}) \Big) \Big]$$

# $AxThExp^{\uparrow}$ :

Inertial observers can move with any speed slower than that of light.



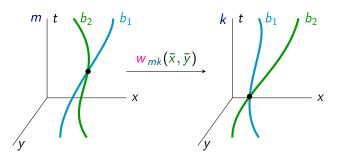
$$\exists h \ [IOb(h)] \land \forall m \bar{x} \bar{y} \ (IOb(m) \land space(\bar{x}, \bar{y}) < c \cdot time(\bar{x}, \bar{y}) \\ \rightarrow \exists k \ [IOb(k) \land W(m, k, \bar{x}) \land W(m, k, \bar{y}) \land m \uparrow k]).$$

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Kin:={AxEField, AxEv, AxSelf, AxSymD, AxLine, AxTriv}
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$$\textbf{NewtonianKin}_{\textit{Full}} := \textbf{Kin} \cup \{ \texttt{AxEther}, \texttt{AbsTime}, \texttt{AxThExp}^{\uparrow}_{+} \}$$

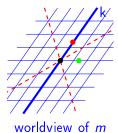
$$\textbf{SpecRel}_{\textit{Full}} := \textbf{Kin} \cup \{ \texttt{AxPh}_c, \texttt{AxThExp}^{\uparrow} \}$$

#### Worldview transformation:



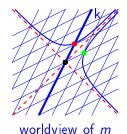
$$\mathbf{w}_{mk}(\bar{x}, \bar{y}) \iff ev_m(\bar{x}) = ev_k(\bar{y})$$

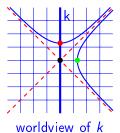
# Galilean transformations:



worldview of k

# Poincaré transformations:





Representation Theorems:

### Theorem:

**SpecRel**<sub>Full</sub>  $\vdash \forall mk[lOb(m) \land lOb(k) \rightarrow "w_{mk} \text{ is a Poincaré}$  Transformation"].

- ca. 1998 proven for a version of BasAx strongly related to SpecRel in the "Big Book" by H. Andréka, J. X. Madarász & I. Németi
- Synthese 2012 "A logic road from special relativity to general relativity" by H. Andréka, J. X. Madarász, I. Németi & G. Székely (Theorem 2.2)

## Theorem:

**NewtonianKin**<sub>Full</sub>  $\vdash \forall mk[IOb(m) \land IOb(k) \rightarrow "w_{mk} \text{ is a Galilean Transformation"}].$ 

 proof for a different version of Newtonian Kinematics in the "Big Book" by H. Andréka, J. X. Madarász & J. Németi

#### Theorem:

**SpecRel**<sub>Full</sub>  $\vdash \forall k P[IOb(k) \land "P \text{ is a (orthochronous) Poincaré}$ Transformation"  $\rightarrow \exists k' w_{kk'} = P].$ 

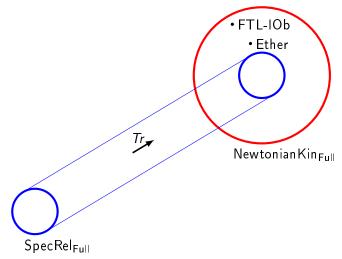
## Theorem:

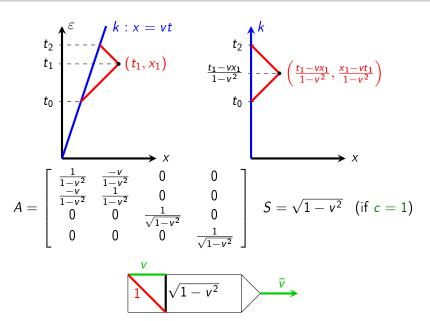
**NewtonianKin**<sub>Full</sub>  $\vdash \forall k G[IOb(k) \land "G \text{ is a (orthochronous)}$ Galilean Transformation"  $\rightarrow \exists k'w_{kk'} = G$ ].

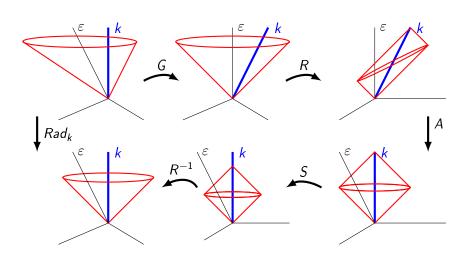
can be proven based on the ideas in the "Big Book" by H.
 Andréka, J. X. Madarász & I. Németi

# Theorem:

There is an interpretation Tr of SpecRel<sub>Full</sub> in NewtonianKin<sub>Full</sub>.







$$Rad_k = R^{-1} \circ S \circ A \circ R \circ G$$



$$Tr(a+b=c) := a+b=c$$
  
 $Tr(a \cdot b = c) := a \cdot b = c$   
 $Tr(a < b) := a < b$ 

$$Tr(\mathbb{W}_{SR}(k,b,t,x,y,z)) := \\ \exists t'x'y'z'[\mathbb{W}_{NK}(k,b,t',x',y',z') \land Rad_k(t',x',y',z') = (t,x,y,z)] \\ Tr(\mathsf{IOb}_{SR}(k)) := \mathsf{IOb}_{NK}(k) \land \forall \varepsilon [Ether(\varepsilon) \rightarrow speed_{\varepsilon}^{NK}(k) < c] \\ \text{where} \\ Ether(\varepsilon) \stackrel{def}{\Longleftrightarrow} \mathsf{IOb}_{NK}(\varepsilon) \land \forall \rho [Ph(\rho) \rightarrow speed_{\varepsilon}^{NK}(\rho) = c]$$

$$\begin{split} \textit{speed}_{\textit{m}}(\textit{k}) = \textit{v} & \iff \forall \bar{\textit{x}}, \bar{\textit{y}} \in \mathsf{Q}^{\textit{d}}[\mathsf{W}(\textit{m},\textit{k},\bar{\textit{x}}) \land \mathsf{W}(\textit{m},\textit{k},\bar{\textit{y}}) \land \textit{x}_{1} \neq \textit{y}_{1} \\ & \rightarrow \frac{\textit{space}(\bar{\textit{x}},\bar{\textit{y}})}{\mathsf{time}(\bar{\textit{x}},\bar{\textit{y}})} = \textit{v}] \end{split}$$

and



$$Tr[\forall m \mid \mathsf{Ob}_{SR}(m) \to \forall \bar{x} \; (\mathsf{W}_{SR}(m, m, \bar{x}) \leftrightarrow x_2 = \ldots = x_d = 0)]$$

$$\equiv \forall m \; [\mathsf{IOb}_{NK}(m) \land \forall \varepsilon [\mathsf{Ether}(\varepsilon) \to \mathsf{speed}_{\varepsilon}^{NK}(m) < c]$$

$$\to \forall \bar{x} \; (\exists \bar{y} [\mathsf{W}_{NK}(m, m, \bar{y}) \land \mathsf{Rad}_k(\bar{y}) = \bar{x}] \leftrightarrow x_2 = \ldots = x_d = 0)]$$

# There is an interpretation Tr of $SpecRel_{Full}$ in $NewtonianKin_{Full}$ . Newtonian $Kin_{Full} \vdash Tr(SpecRel_{Full})$

- AxEField<sub>NK</sub> ⊢ Tr(AxEField<sub>SR</sub>)
- $AxEField_{NK}$ ,  $AxLine_{NK}$ ,  $AxEther_{NK}$ ,  $AxEv_{NK} \vdash Tr(AxEv_{SR})$
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub>, AxSelf<sub>NK</sub> ⊢ Tr(AxSelf<sub>SR</sub>)
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub>, AxSymD<sub>NK</sub> ⊢
   Tr(AxSymD<sub>SR</sub>)
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub> ⊢ Tr(AxLine<sub>SR</sub>)
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub>, AxTriv<sub>NK</sub> ⊢ Tr(AxTriv<sub>SR</sub>)
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub> ⊢ Tr(AxPh<sub>cSR</sub>)
- AxEField<sub>NK</sub>, AxLine<sub>NK</sub>, AxEther<sub>NK</sub>, AxThExp<sup>↑</sup><sub>+NK</sub> ⊢
   Tr(AxThExp<sup>↑</sup><sub>SP</sub>)



# $AxThExp^{\uparrow}_{NoFTL}$ :

Inertial observers can move with any speed which is in the ether frame slower than that of light.

$$\exists h \ (IOb(h)) \land \forall \varepsilon \overline{x} \overline{y} \big( Ether(\varepsilon) \land \operatorname{space}(\overline{x}, \overline{y}) < c \cdot \operatorname{time}(\overline{x}, \overline{y}) \\ \rightarrow \exists k \ IOb(k) \land W(\varepsilon, k, \overline{x}) \land W(\varepsilon, k, \overline{y}) \land m \uparrow k \big).$$

### AxNoFTL:

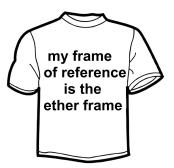
All inertial observers move slower than light with respect to the ether frames.

$$\forall m \varepsilon [IOb(m) \land Ether(\varepsilon) \rightarrow Speed_{\varepsilon}^{NK}(m) < c].$$

# $SpecRel_{Full}^{\varepsilon} := SpecRel_{Full} \cup \{AxPrimitiveEther\}$

### AxPrimitiveEther:

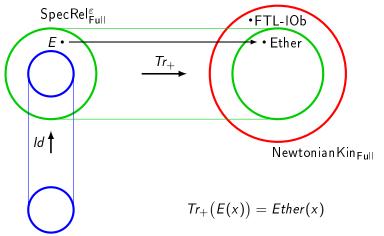
There is a non-empty class of ether observers, stationary with respect to each other, which is closed under trivial transformations.



$$\exists \varepsilon \big( E(\varepsilon) \land \forall k \big[ [IOb(k) \land (\exists T \in Triv) w_{\varepsilon k}^{SR} = T] \leftrightarrow E(k) \big] \big)$$



 $Tr_+$  is a definitional equivalence between  $SpecRel_{Full}^{\varepsilon}$  and NewtonianKin NoFTL.



# $Tr_+$ is a definitional equivalence between $\mathbf{SpecRel}_{Full}^{\varepsilon}$ and $\mathbf{NewtonianKin}_{Full}^{NoFTL}$ .

